



**Charm elliptic flow
from
parton coalescence dynamics**

Denes Molnar

Ohio State University, Columbus, OH, USA

Hot Quarks 2004

July 18-25, Taos Ski Valley, NM

The plan

- Motivation

- (Rafting) theory

- parton coalescence + transport theory

- Results (on/in the water)

- parton chemistry & elliptic flow

- charm meson flow

- a few puzzles

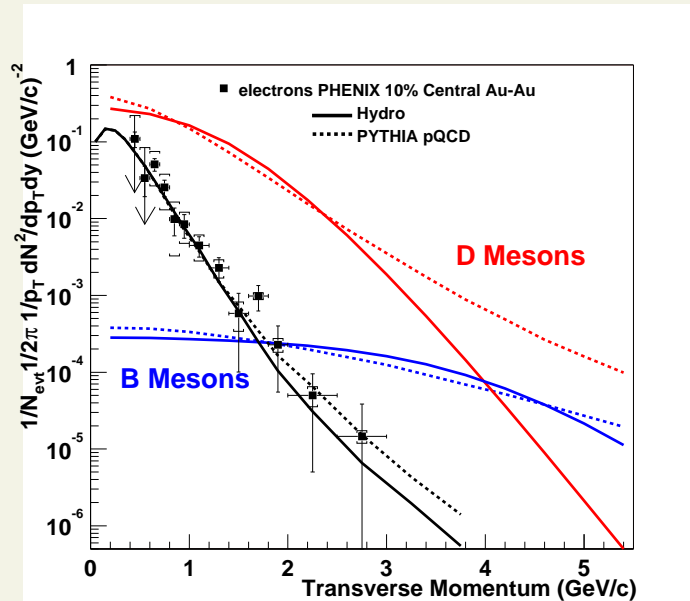
Why charm?

- **additional, heavy, probe** of the same partonic medium
 - perturbative QCD: expect **weaker jet-quenching for heavy quarks**

[Djordjevic, Gyulassy ('03), Dokshitzer, Kharzeev ('01)]

- **charm thermalization or purely jet production?**

- **both scenarios seem to be consistent with PHENIX electron spectra** [Batsouli et al ('03)]



→ **charm quark elliptic flow is the key**: tells amount of rescattering

- **Also: additional test for quark coalescence models** (flow addition formulas)

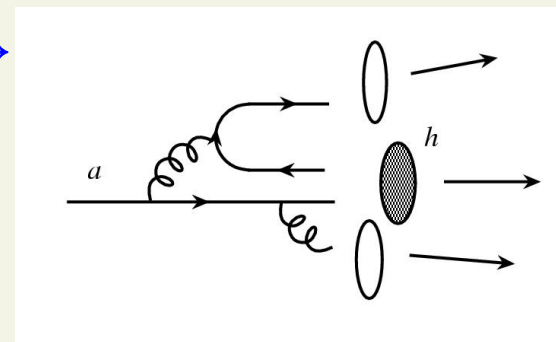
First some theory:

- Parton coalescence
- Covariant parton transport

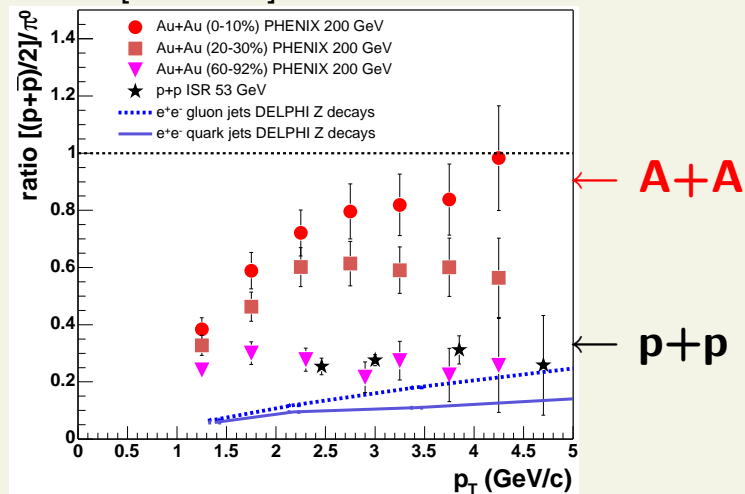


Parton coalescence

- independent fragmentation: $D_{a \rightarrow h}(z) \rightarrow$
gives $p/\pi \approx 0.2 \Rightarrow$ fails for baryons at RHIC

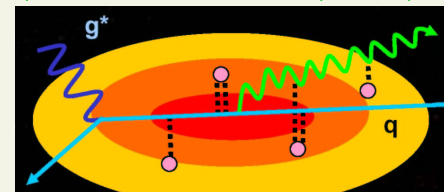


Esumi [PHENIX]:



problem persists for jet quenching:

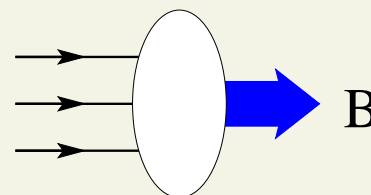
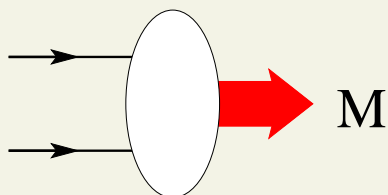
Wang, Gyulassy, Dokshitzer, Mueller, Levai, Vitev, Wiedemann et al, Guo, Djordjevic, ...



- parton coalescence

Hwa, Yang, Biró, Zimányi, Lévai, Csizmadia, Ko, Lin, Voloshin, D.M., Greco, Fries, Müller, Nonaka, Bass, ...

heavy-ion collisions - large parton density \Rightarrow multi-parton processes also relevant



lowest-order $q\bar{q} \rightarrow M, qqq \rightarrow B$ (valence quarks only)

Applications

- hadron multiplicity:** Das & Hwa, PLB68 ('77)
Biró et al, PLB347 ('95) - ALCOR
Csizmadia & Lévai, JPG28 ('02) - MICOR
- baryon/meson ratio:** Hwa & Yang, PRC65 ('02)
Greco, Ko, Levai, PRL90 ('03); PRC68 ('03)
Fries, Müller, Nonaka, Bass, PRL90 ('03); PRC68 ('03)
Hwa, Yang, PRC67 ('03)
Fries, Müller, Nonaka, Bass, JPG30 ('04)
Hwa & Yang, nucl-th/0401001
- resonances:** Nonaka, Müller, Asakawa, Bass, Fries, PRC69 ('04),
Zimányi & Lévai, nucl-th/0404060
- elliptic flow:** Ko & Lin, PRL89 ('02)
Voloshin, NPA715 ('02)
D.M. & Voloshin, PRL91 ('03)
Nonaka, Fries, Bass, PLB583 ('04)
D.M., JPG30 ('04)
Greco, Ko, nucl-th/0404020
D.M., nucl-th/0403035
- charm hadron elliptic flow:** Lin & D.M., PRC68 ('03)
Greco, Ko, Rapp, nucl-th/0312100

Simple coalescence formula

- based on: $n + p \rightarrow d$

Butler & Pearson, PR129 ('63); Schwarzschild & Zupancic, PR129 ('63); Sato & Yazaki, PLB98 ('81); Gyulassy, Frankel & Remler, NPA402 ('86); Dover et al PRC44 ('91); Nagle et al PRC53 ('96) Kahana et al PRC54 ('96); Scheibl & Heinz, PRC59 ('99); ...

• basic equations: $q\bar{q} \rightarrow meson$, qqq (or $\bar{q}\bar{q}\bar{q}$) $\rightarrow baryon$

$$\frac{dN_M(\vec{p})}{d^3p} = g_M \int \left(\prod_{i=1,2} d^3x_i d^3p_i \right) W_M(x_1 - x_2, \vec{p}_1 - \vec{p}_2) f_\alpha(\vec{p}_1, x_1) f_\beta(\vec{p}_2, x_2) \delta^3(\vec{p} - \vec{p}_1 - \vec{p}_2)$$

$$\frac{dN_B(\vec{p})}{d^3p} = g_B \int \left(\prod_{i=1,2,3} d^3x_i d^3p_i \right) W_B(x_{12}, x_{13}, \vec{p}_{12}, \vec{p}_{13}) f_\alpha(\vec{p}_1, x_1) f_\beta(\vec{p}_2, x_2) f_\gamma(\vec{p}_3, x_3) \delta^3(\vec{p} - \sum \vec{p}_i)$$

hadron yield space-time hadron wave-fn. quark distributions

\approx “partons close in phasespace can coalesce”

assumes: weak binding, no 2-body or 3-body correlations, rare process,
sudden hadronization

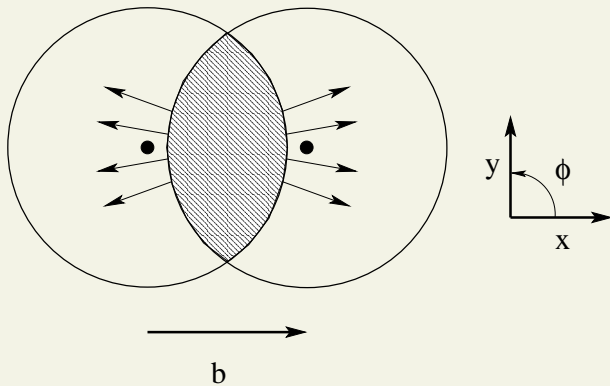
few \rightarrow many mapping: 8 f_q 's (u, d, s, c) give spectra for all hadron species

Two approaches

- **Parameterize f_q :** \sim spirit of HBT source params, or Blastwave
 - this is what have been mostly done so far

successes:

- **baryon/meson ratio enhancement at RHIC** [Greco et al, Fries et al, Hwa et al, ...]
- **elliptic flow (v_2) scaling with quark number** [Voloshin, D.M., Lin,...]



$$v_2^M(p_\perp) \approx v_2^a(z_a p_\perp) + v_2^{\bar{b}}(z_b p_\perp)$$

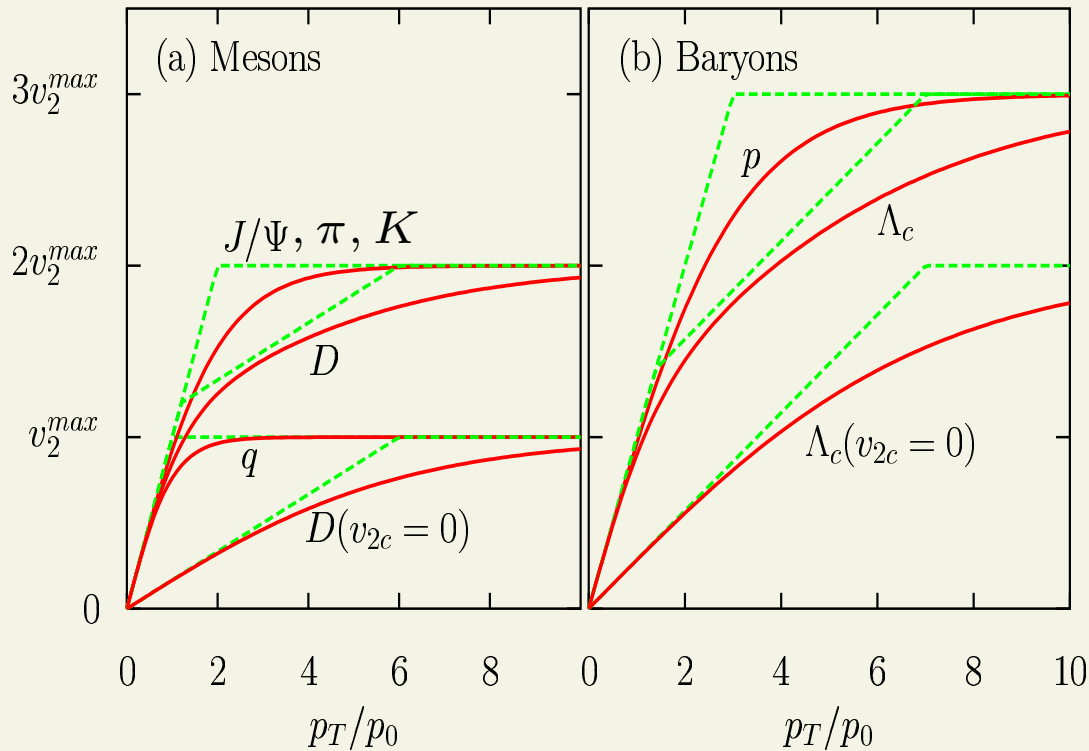
$$v_2^B(p_\perp) \approx v_2^a(z_a p_\perp) + v_2^b(z_b p_\perp) + v_2^c(z_c p_\perp)$$

$$z_a : z_b[: z_c] \approx m_a : m_b[: m_c]$$

- **Compute f_q :** the only way to make predictions
 - requires a **dynamical** model
 - here choose: **covariant transport theory**

Coalescence expectations for charm

Lin & D.M. ('03)



$$m_q : m_s : m_c = 3 : 5 : 15$$

$$v_2^D(p_\perp) \approx v_2^c\left(\frac{5p_\perp}{6}\right) + v_2^q\left(\frac{p_\perp}{6}\right)$$

$$v_2^{D_s}(p_\perp) \approx v_2^c\left(\frac{3p_\perp}{4}\right) + v_2^q\left(\frac{p_\perp}{4}\right)$$

$$v_2^{\Lambda_c}(p_\perp) \approx v_2^c\left(\frac{5p_\perp}{7}\right) + 2v_2^q\left(\frac{p_\perp}{7}\right)$$

$$v_2^{J/\Psi}(p_\perp) \approx 2v_2^c\left(\frac{p_\perp}{2}\right)$$

simplicistic (linear & flat); more realistic (based on parton cascade MPC)

- $v_2(p_\perp)$ rises slower, saturates later for asymmetric systems D, D_s, Λ_c
 - heavy quark carries most of hadron momentum (momentum \propto constituent mass)
- nonzero v_2 for D, D_s, Λ_c , even for zero charm v_2 (no thermalization)

Covariant transport theory

[Pang, Zhang, Gyulassy, D.M., Vance, Csizmadia, Pratt, Cheng, ...]

- **Transport equation:** - kinetic theory of a parton gas \sim Boltzmann

$$p^\mu \partial_\mu f_i(x, \vec{p}) = S_i + \sum_{jkl} C_{ij \rightarrow kl}^{2 \rightarrow 2} [f_i, f_j, f_k, f_l] + \dots C^{2 \leftrightarrow 3} \dots + \dots$$

f_i : 1-particle phasespace distr., $C^{n \leftrightarrow m}$: contains the $n \rightarrow m$ matrix elements

S_i : initial conditions

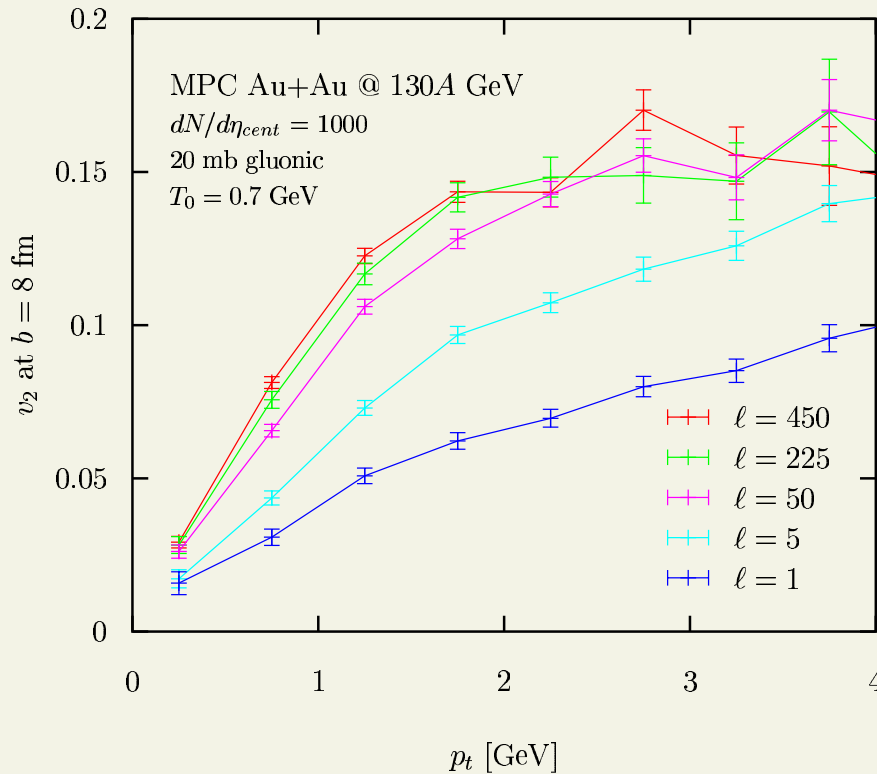
- relevant parameter is the **mean free path** $\lambda_{MFP} \sim 1/\rho\sigma$
- interpolates between ideal hydro $\lambda = 0$ and free streaming $\lambda = \infty$

- **solvable only numerically:** **nonlinear integro-diff eq. in 6D phasespace**
 - few Lorentz-covariant algorithms: **ZPC**, **MPC**, PSYCHE, GROMIT
 - common element: **parton subdivision** - essential to maintain covariance

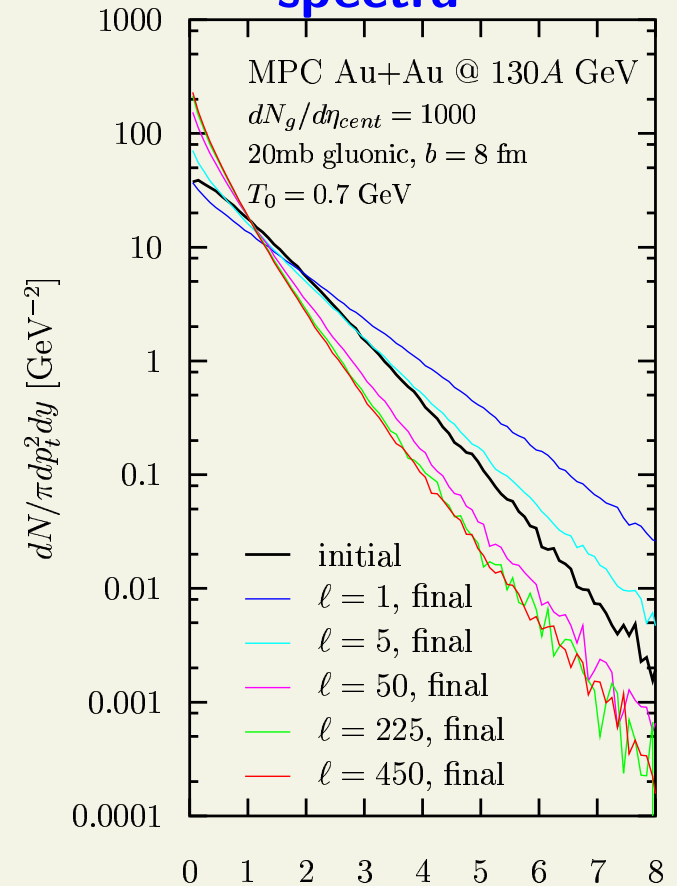
Subdivision matters

rescaling trick $f \rightarrow f \cdot l, \sigma \rightarrow \sigma/l$ - $l \rightarrow \infty$ gives covariant answer

D.M. & Gyulassy ('02): $v_2(p_T)$



spectra



- naive cascade ($l = 1$): numerical artifacts - smaller v_2 , hotter spectra
- need subdivision $l \sim 200$

Transport + coalescence

Combine **parton transport** (\rightarrow charm quark v_2) & **coalescence** (\rightarrow charm hadron v_2)

- run parton cascade until freezeout (g, d, u, s, c) (MPC)
- then do parton coalescence (Gyulassy-Frankel-Remler formula, NPA 402 ('83))
- hadronize partons w/o coalescence partner via indep. fragmentation (JETSET)

Processes: leading-order perturbative QCD

elastic: $gg \rightarrow gg, gq \rightarrow gq, qq \rightarrow qq, q\bar{q} \rightarrow q\bar{q}, \bar{q}\bar{q} \rightarrow \bar{q}\bar{q}$

inelastic: $gg \leftrightarrow q\bar{q}, q\bar{q} \rightarrow q'\bar{q}'$

Initial conditions: Au+Au at $\sqrt{s} = 200A$ GeV, $b = 8$ fm:

$p_T > 2$ GeV: **minijets** from leading-order pQCD

$p_T < 2$ GeV: soft component via **smoothly extrapolated spectra**

normalized to yield $dN_{parton}(b=0)/dy = 2000$

Processes

- drive calculation with $\sigma_{gg \rightarrow gg} = 10 \text{ mb}$
- relative “strength” of channels:

elastic:

$$gg \rightarrow gg \equiv 1 \text{ unit}$$

$$gq \rightarrow gq = (4/9) \approx 0.5,$$

$$qq \rightarrow qq = (4/9)^2 \approx 0.2$$

roughly same for $q = c$

inelastic:

$$gg \rightarrow q\bar{q} \leq 1/27 \approx 0.04,$$

$$q\bar{q} \rightarrow gg \leq (16/6)^2 \times 1/27 \approx 0.3$$

$$q\bar{q} \rightarrow q'\bar{q}' \leq 16/243 \approx 0.07$$

$$gg \rightarrow c\bar{c} \approx 0.07\mu_D^2/M_c^2 \approx 0.02,$$

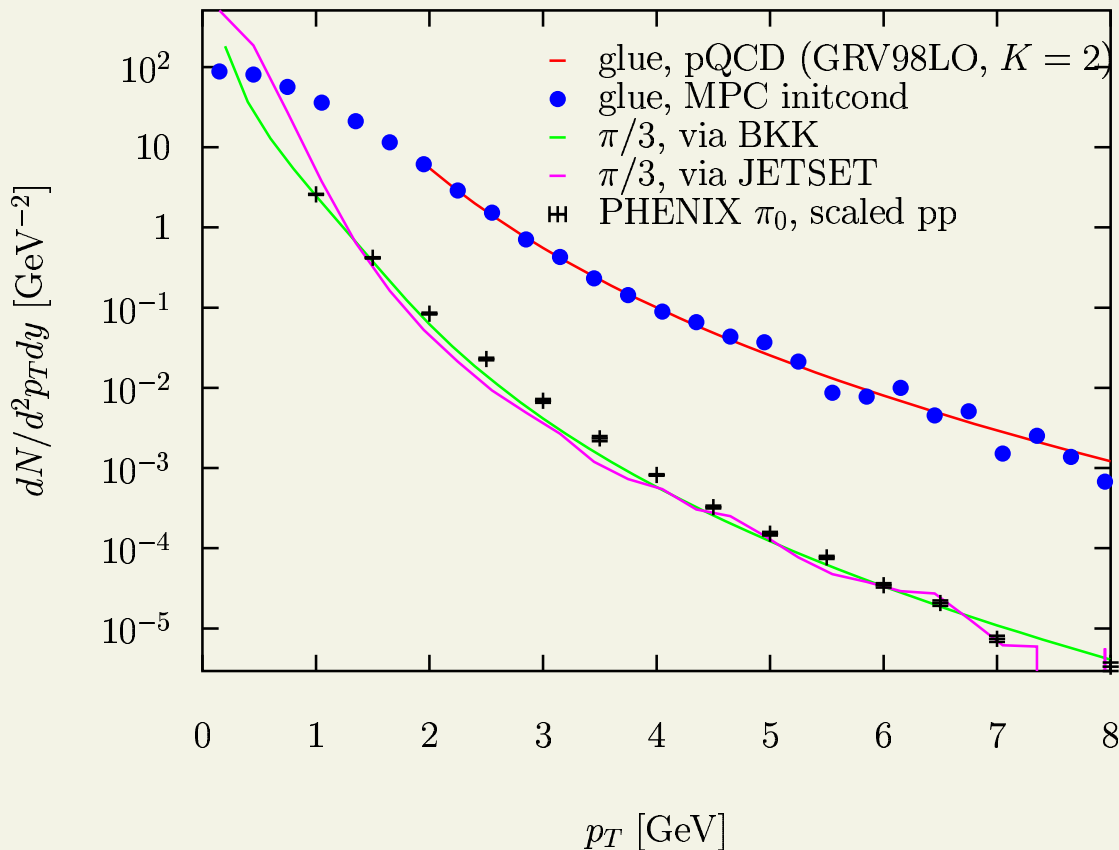
$$c\bar{c} \rightarrow gg \approx 0.5\mu_D^2/M_c^2 \approx 0.2$$

$$q\bar{q} \rightarrow c\bar{c} \leq 0.009\mu_D^2/M_c^2 \approx 0.003,$$

$$c\bar{c} \rightarrow q\bar{q} \leq 0.009\mu_D^2/M_c^2 \approx 0.003$$

Initial conditions

- Au+Au at RHIC with $b = 8$ fm, i.e., 30% centrality
- $p_{\perp} > 2$ GeV: minijets(dijets) [GRV98LO, $K = 2$]
 $p_{\perp} < 2$ GeV: smoothly joined-on soft component, such that $dN^{parton}/dy(b = 0) = 2000$
- binary collision transverse profile, formation time $\tau_0 = 0.1$ fm/c



Ready to paddle! - results

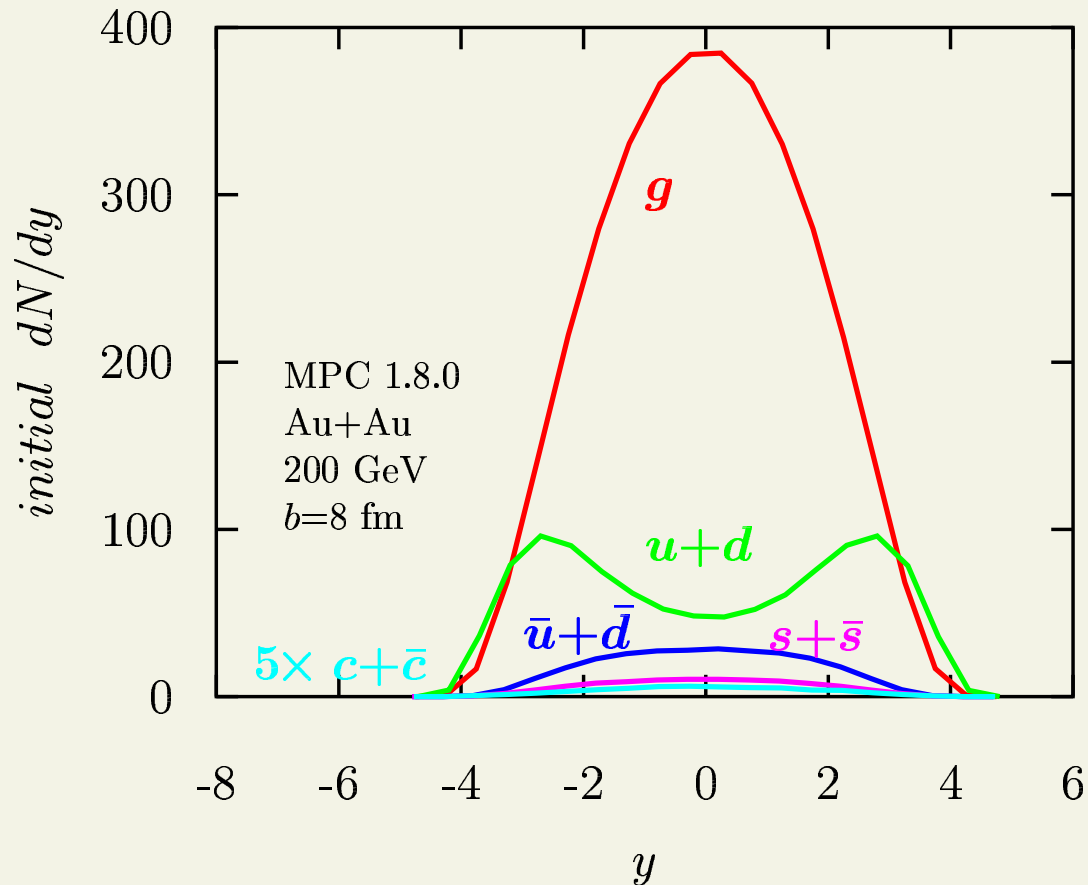
- Chemistry - secondary charm production
- Charm quark elliptic flow
- Prediction for D mesons



Charm production

D.M. ('04):

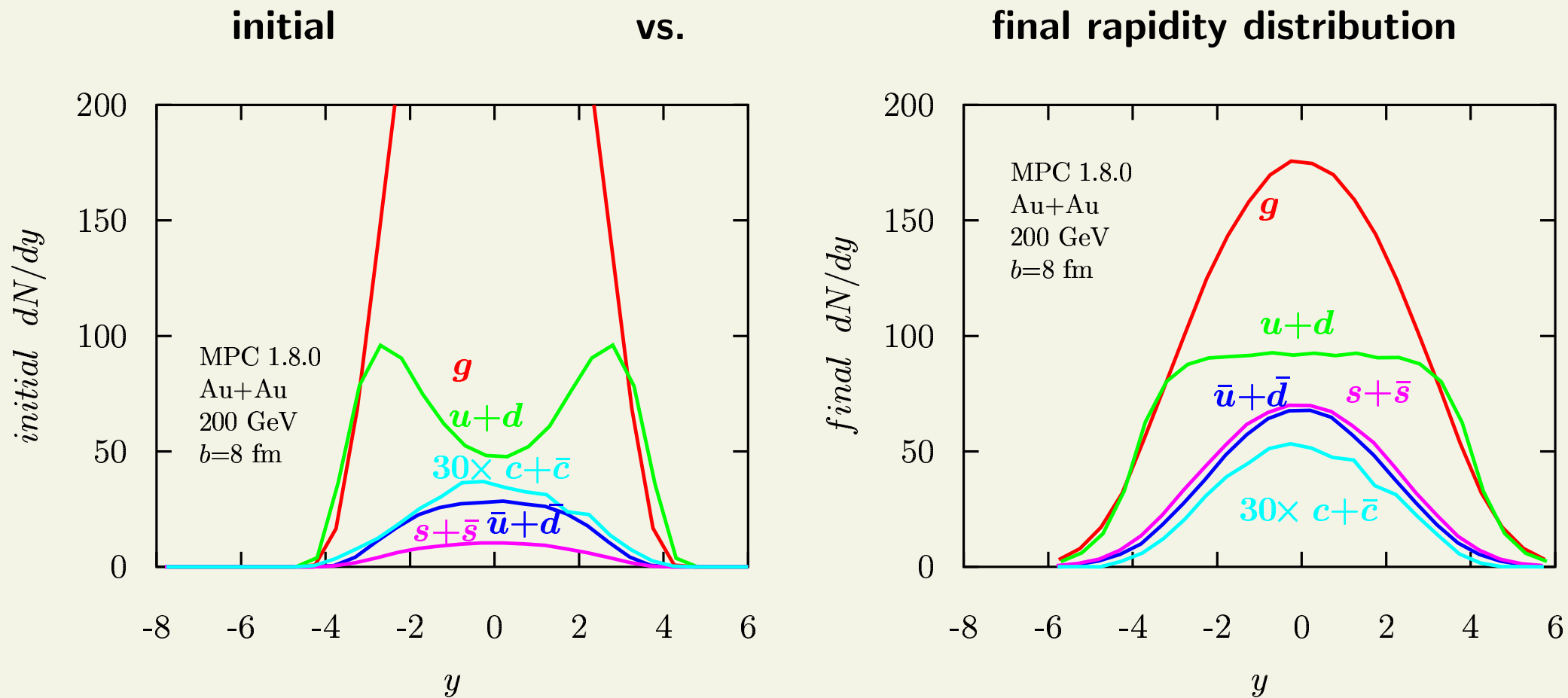
initial rapidity distributions



- **gluons dominate at RHIC**

Charm production

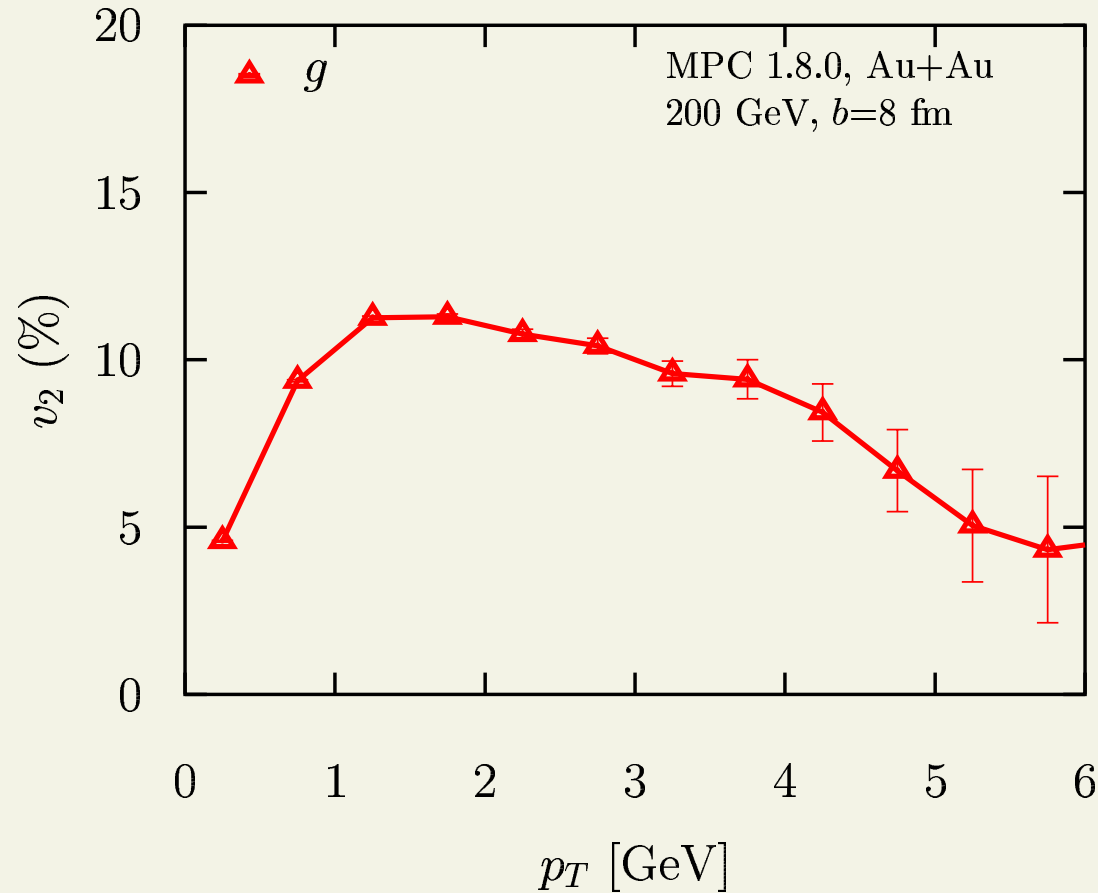
D.M. ('04):



- roughly half the glue fuse to $q\bar{q}$
- extra 40 – 50% charm yield due to secondary production
- strangeness is up by much more, factor 5 or so

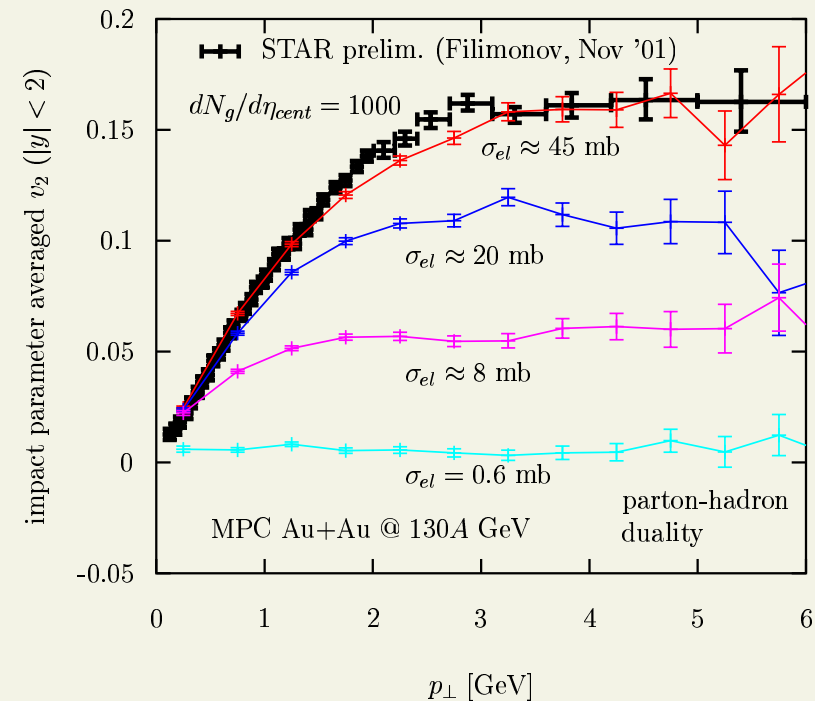
Charm quark elliptic flow

D.M. ('04):



D.M. & Gyulassy ('02):

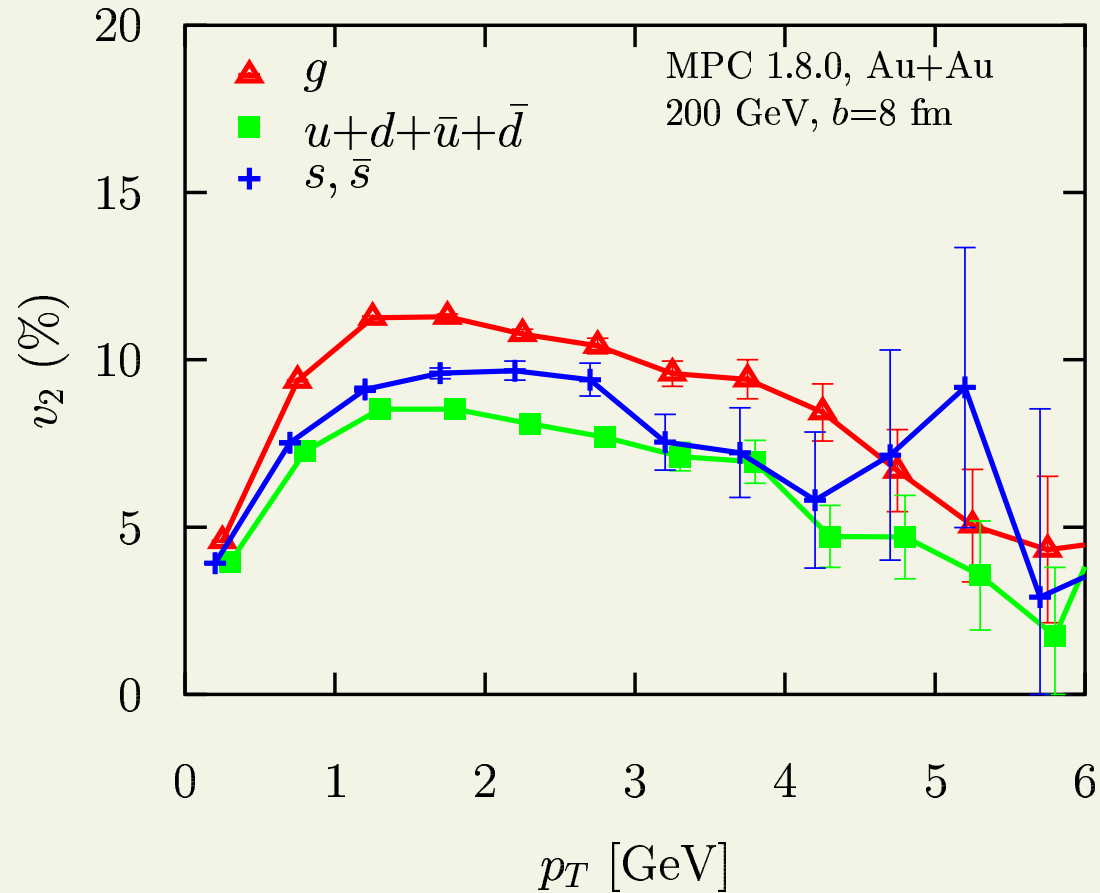
generic earlier results



- **turnover and decrease** at high p_T seen for the first time
 - difference in initial conditions?

Charm quark elliptic flow

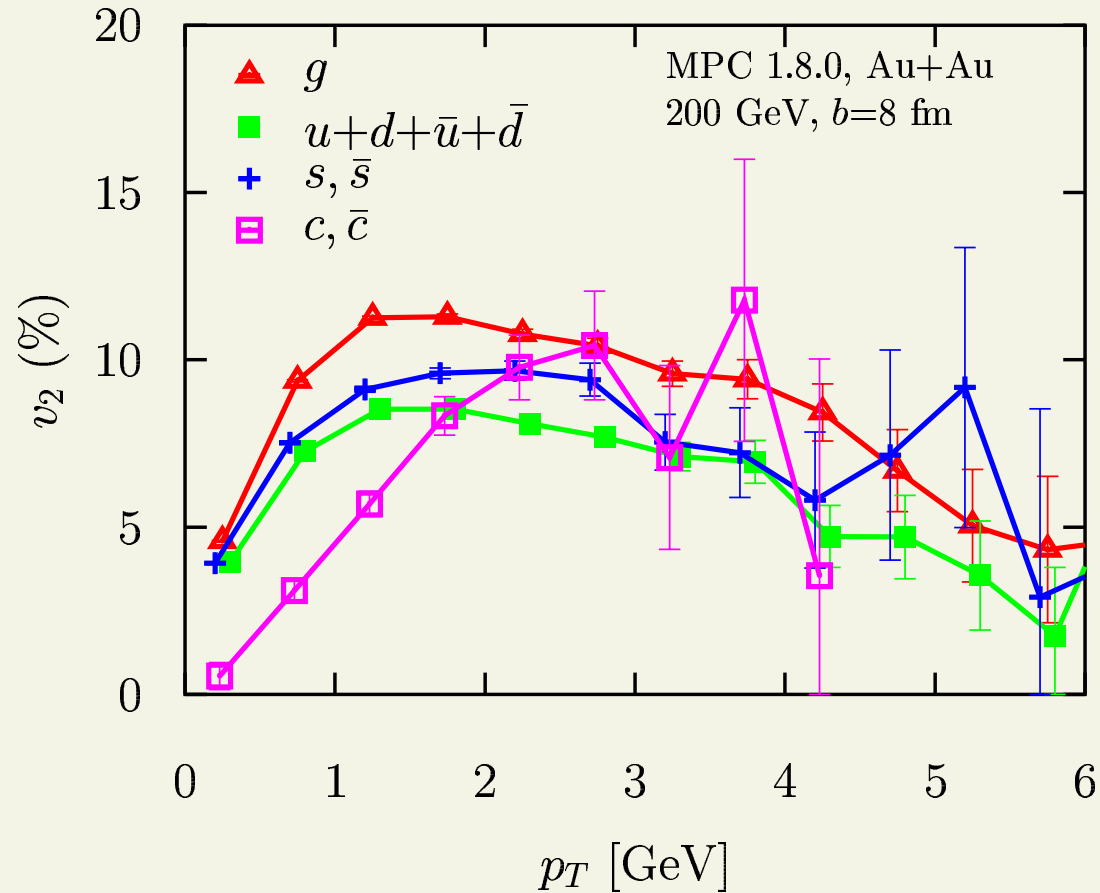
D.M. ('04):



- **light quark v_2 's are smaller** because mainly driven by $gq \rightarrow gq$
 - 4/9 times weaker than $gg \rightarrow gg$

Charm quark elliptic flow

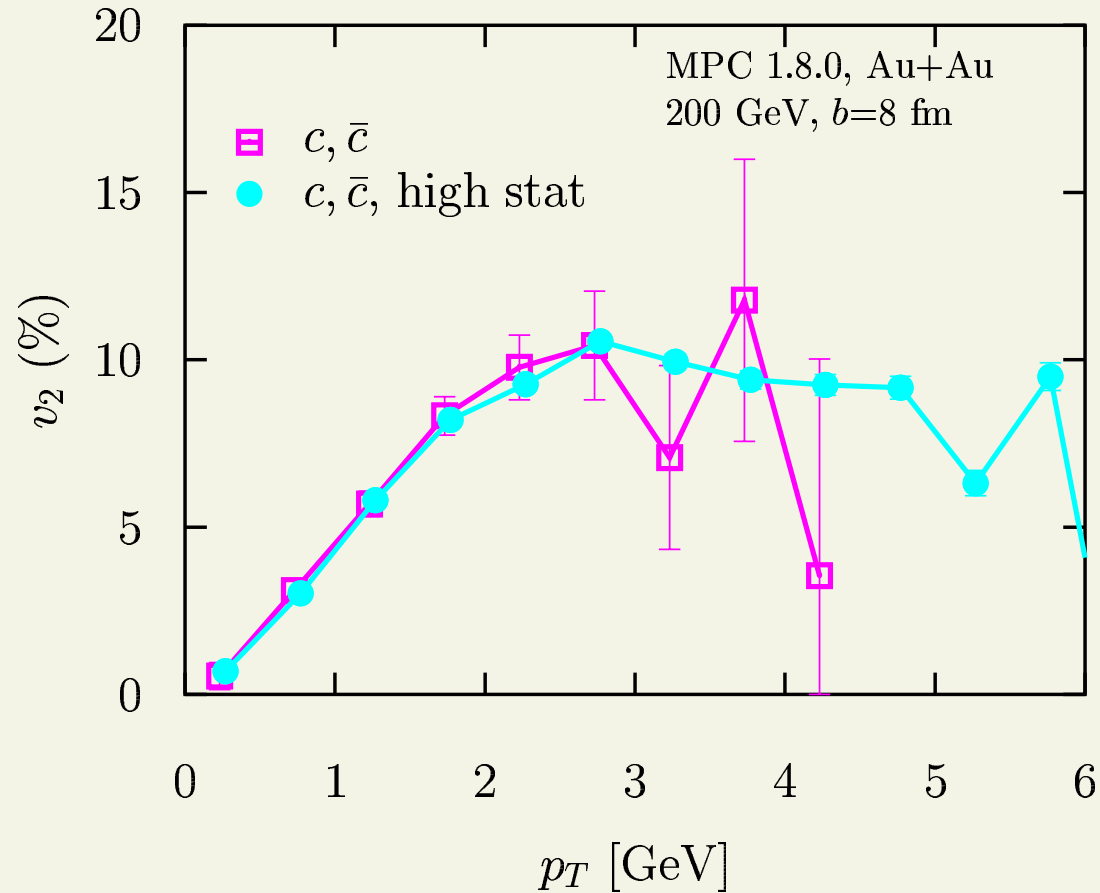
D.M. ('04):



- charm v_2 rises slower at low $p_T \sim$ hydro mass effect
- (!) but becomes comparable to light quark v_2 above $p_T \approx 2.5$ GeV (!)

Charm quark elliptic flow

D.M. ('04):

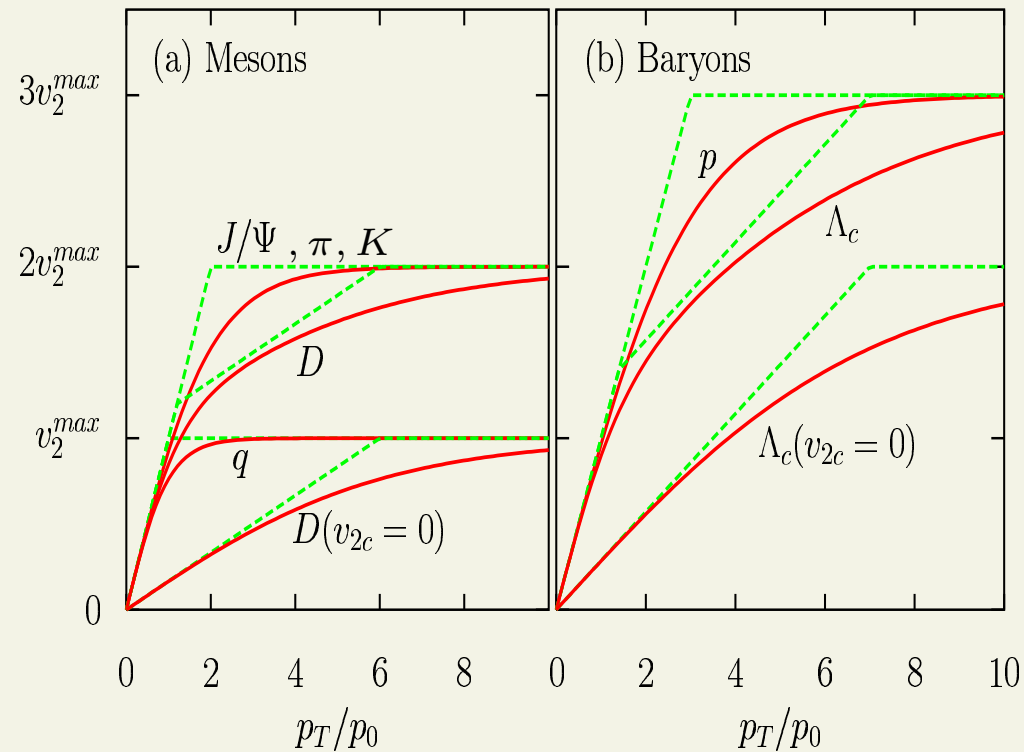
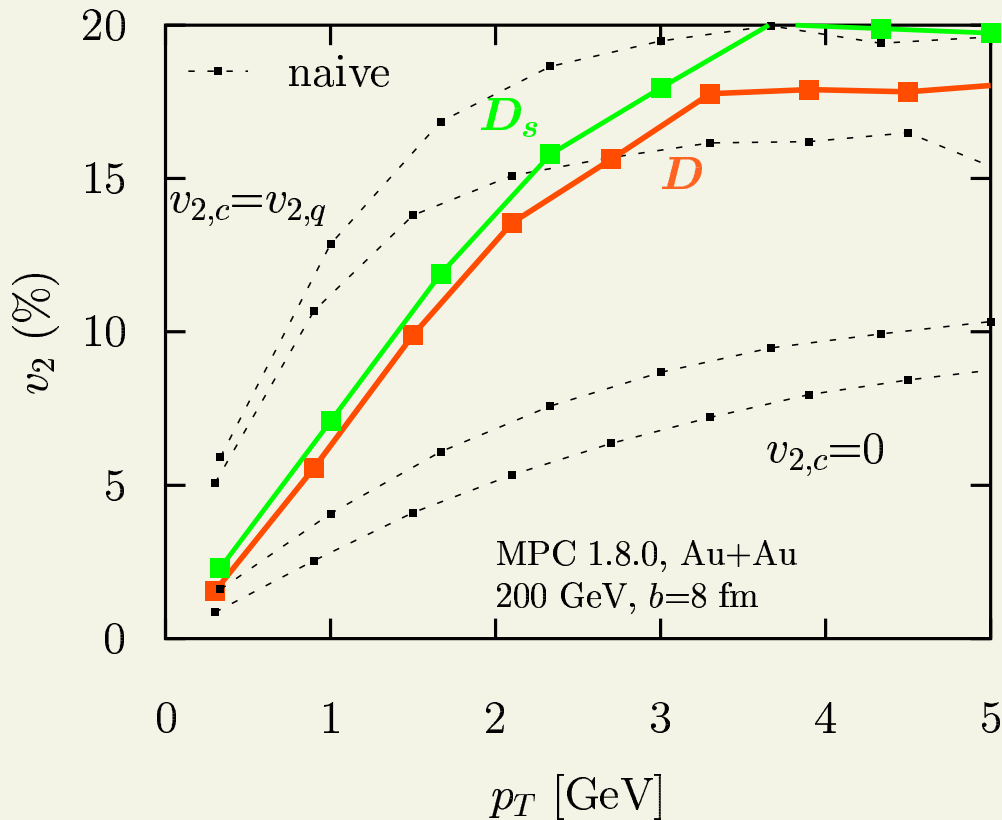


- enhanced high- p_T stat. using “perturbative” charm (no feedback on medium)
- consistent with equal-weight sampling

D meson v_2 from coalescence

D.M. ('04):

[naive: Lin & D.M. ('03)]



- for D , **later saturation ($p_T \sim 3$ GeV)** than for light mesons ($p_T \sim 2$ GeV)
- very different from scenario w/o charm flow
- however, eventually similar v_2 magnitude - twice that of quarks

Not all stunts work - open issues

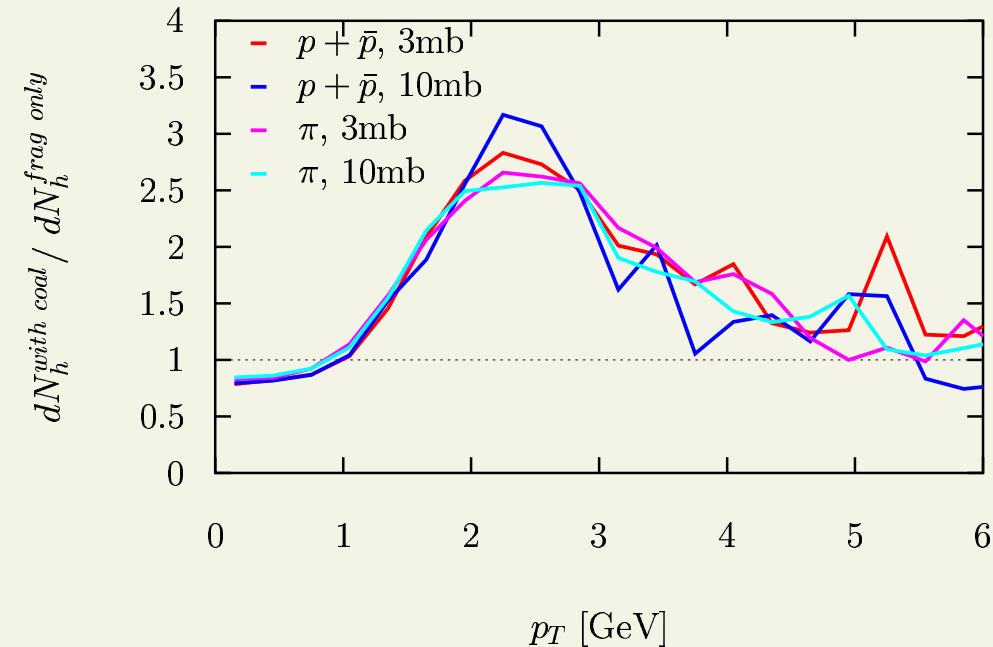
- Fragmentation contribution



Puzzles for light quarks

B/M unchanged, same enhancement for π & p

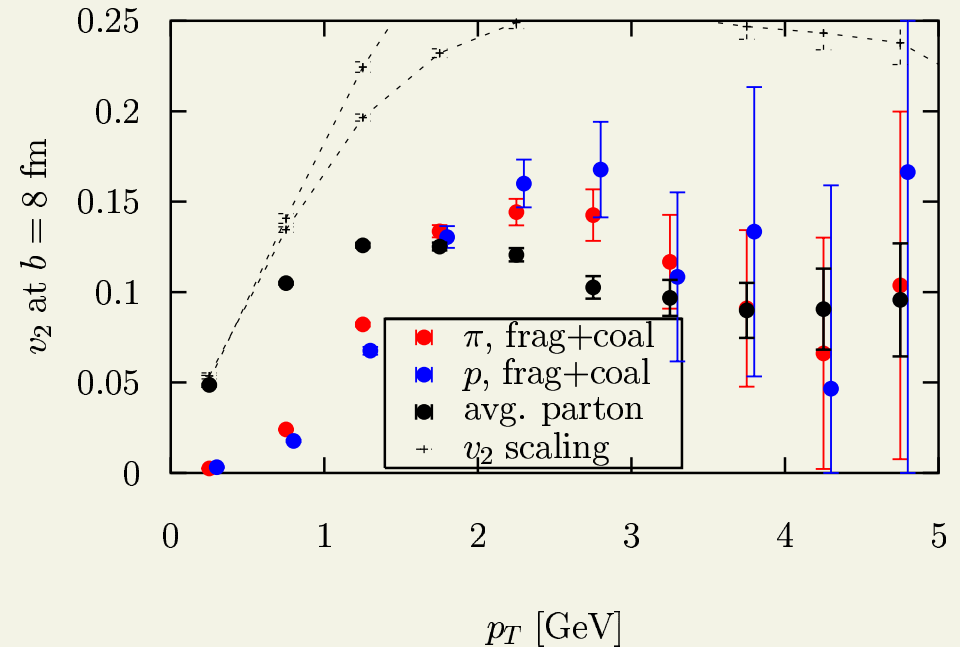
D.M., nucl-th/0403035 ('04):



- $2 - 3\times$ **enhancement** over fragmentation for $1.5 < p_T < 4$ **GeV**

but same for protons as pions $\Rightarrow p/\pi$ stays low
 \rightarrow **baryons have more constituents** \Rightarrow are more "fragile"

frag. contribution suppresses v_2

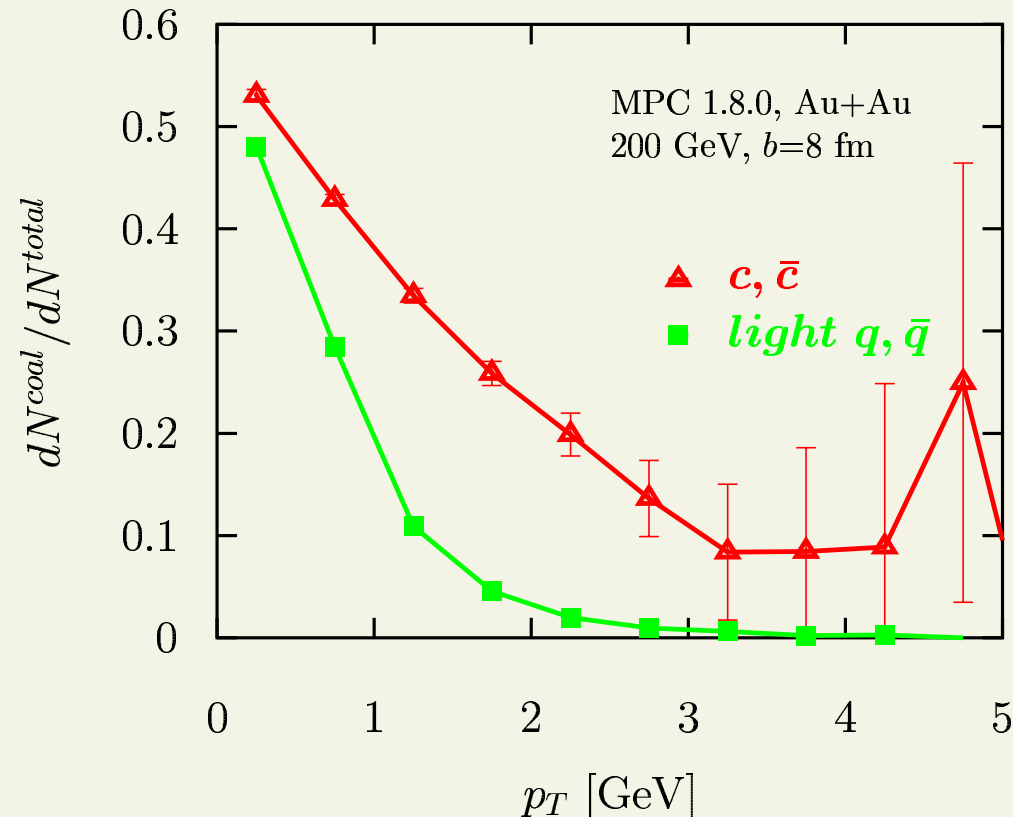


- **still 25-30% of partons fragment**
 \rightarrow reduced v_2 relative to pure coalescence
- **baryon-meson splitting almost gone**

Need higher freezeout density (smaller σ) - but how to keep then v_2 large??

Same puzzle for charm

D.M. ('04): **key issue - fraction of partons that coalesce**



- **80-90% of charm quarks above $p_T > 3$ GeV fragment**
 - fragment roughly preserving the quark momentum ($z \approx 1$)
 - for these hadrons elliptic flow is about the same as for charm quarks - **no doubling**
- ⇒ **coalescence v_2 contribution will be reduced above $p_T \sim 3 - 4$ GeV**

Trip is over - remember the best parts

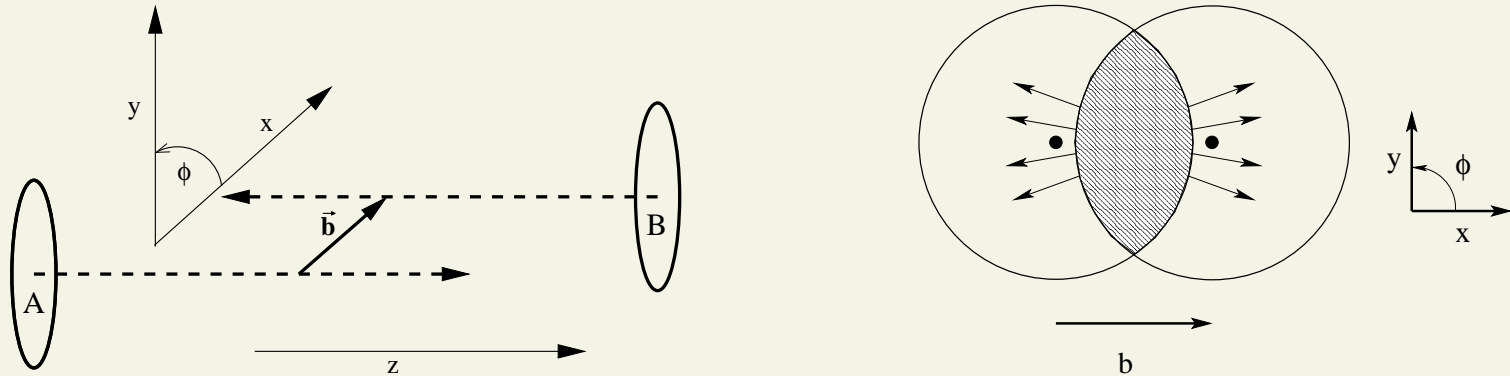
- 40-50% extra secondary charm from transport
- significant charm v_2 at high p_T
- D meson $v_2(p_T)$ may saturate at similar values as that of light mesons, but at 50% higher p_T

- already planning next one:
 - “best chance” scenario for coalescence
 - various other improvements

Additional ammunition

Azimuthal anisotropy (v_2, \dots)

- momentum-space **anisotropy** of particle production in **A+A** collisions



$$\frac{dN}{d\phi dX} \equiv \frac{1}{2\pi} \frac{dN}{dX} \left[1 + 2 \sum_{n=1} v_n(X) \cos(n\phi) \right] \rightarrow \begin{array}{l} \text{directed flow: } v_1(X) \equiv \langle \cos \phi \rangle_X \\ \text{elliptic flow: } v_2(X) \equiv \langle \cos 2\phi \rangle_X \end{array}$$

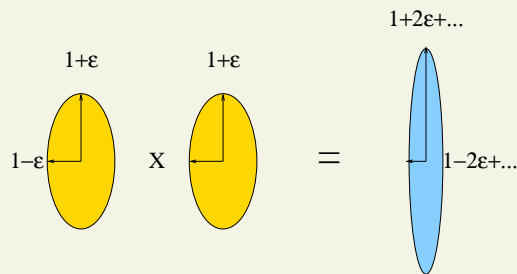
X : event and particle selection, e.g., centrality, **transverse momentum**

- v_2 is a measure of the opacity, or number of scatterings
 - w/o scatterings, $v_2 = 0$

Anisotropy amplification

[D.M. & Voloshin, PRL91 ('03)]

narrow wave fn. (in p): $q\bar{q} \Rightarrow \longrightarrow M \quad (2 \times \frac{\vec{p}}{2} \rightarrow \vec{p}), \quad 3q \Rightarrow \longrightarrow B$
 $(3 \times \frac{\vec{p}}{3} \rightarrow \vec{p})$



$$\frac{dN_M}{d\phi} \propto \left(\frac{dN_q}{d\phi} \right)^2, \quad \frac{dN_B}{d\phi} \propto \left(\frac{dN_q}{d\phi} \right)^3$$

$$v_2^M(p_\perp) \approx v_2^a\left(\frac{p_\perp}{2}\right) + v_2^{\bar{b}}\left(\frac{p_\perp}{2}\right)$$

$$v_2^B(p_\perp) \approx v_2^a\left(\frac{p_\perp}{3}\right) + v_2^b\left(\frac{p_\perp}{3}\right) + v_2^c\left(\frac{p_\perp}{3}\right)$$

⇒ **hadron v_2 amplified at high p_\perp**

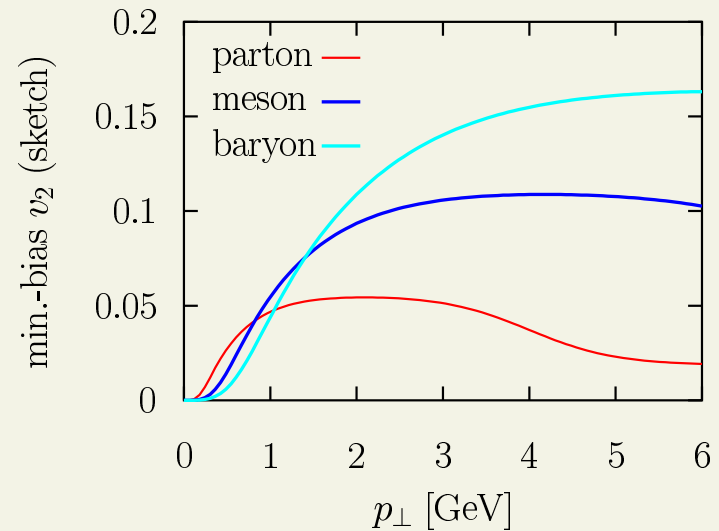
if all quarks have same v_2 :

$3\times$ for baryons } **50% larger v_2**
 $2\times$ for mesons } **for baryons**

$$\text{“ } v_2^h(p_\perp) \approx n \times v_2^q(p_\perp/n) \text{ ”}$$

→ $5\times$ for pentaquark, $6\times$ for deuteron

- if v_2 depends on quark flavor, further differentiation by flavor content



Determination of parton v_2

Via measuring $v_2(p_\perp)$ for various hadrons, one can extract quark flows AND test consistency of coalescence model

Consider, e.g., $v_2^q \neq v_2^s$ (light vs. strange)

$$\begin{array}{ll}
 v_2^\pi(p_\perp) \approx 2v_2^q\left(\frac{p_\perp}{2}\right) & v_2^p(p_\perp) \approx 3v_2^q\left(\frac{p_\perp}{3}\right) \\
 v_2^K(p_\perp) \approx v_2^q\left(\frac{p_\perp}{2}\right) + v_2^s\left(\frac{p_\perp}{2}\right) & v_2^{\Lambda,\Sigma}(p_\perp) \approx 2v_2^q\left(\frac{p_\perp}{3}\right) + v_2^s\left(\frac{p_\perp}{3}\right) \\
 v_2^\phi(p_\perp) \approx 2v_2^s\left(\frac{p_\perp}{2}\right) & v_2^\Xi(p_\perp) \approx v_2^q\left(\frac{p_\perp}{3}\right) + 2v_2^s\left(\frac{p_\perp}{3}\right) \\
 & v_2^\Omega(p_\perp) \approx 3v_2^s\left(\frac{p_\perp}{3}\right)
 \end{array}$$

2 unknowns, 7 equations \Rightarrow e.g.,

$$\begin{aligned}
 v_2^q(p_\perp) &= v_2^\pi(2p_\perp)/2 = v_2^p(3p_\perp)/3 = v_2^\Lambda(3p_\perp) - 2v_2^p(3p_\perp)/3 \\
 v_2^s(p_\perp) &= v_2^\phi(2p_\perp)/2 = v_2^\Omega(3p_\perp)/3 = v_2^K(2p_\perp) - v_2^\pi(2p_\perp)/2
 \end{aligned}$$

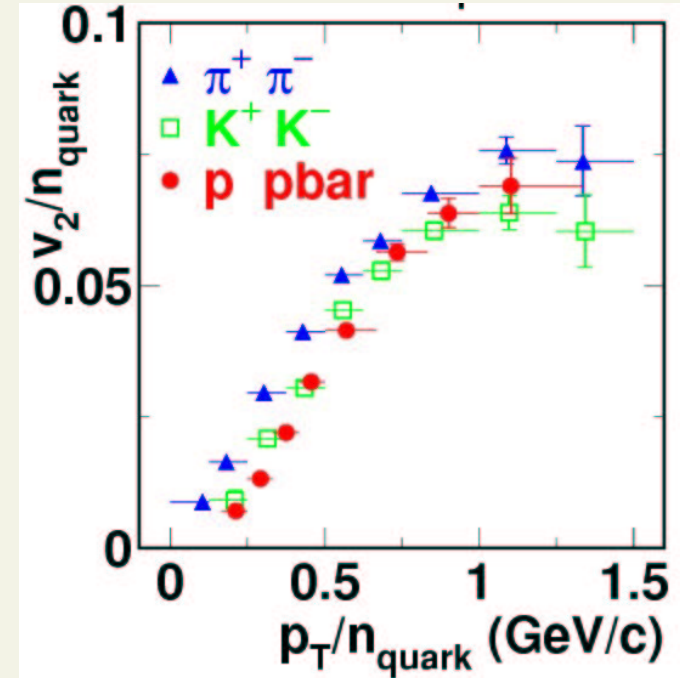
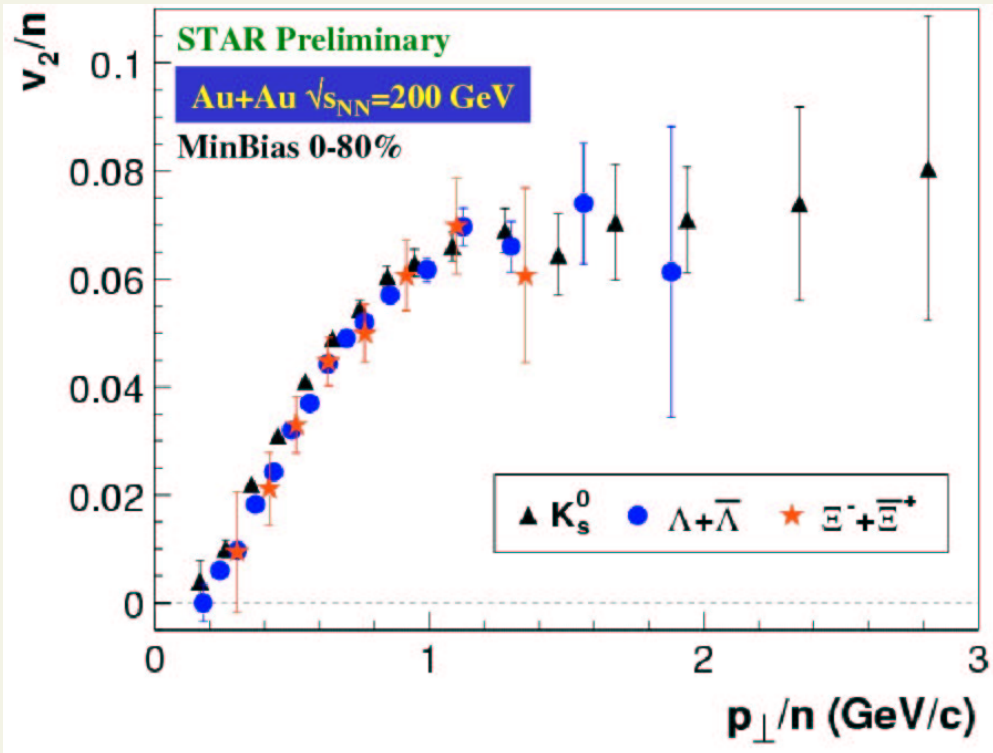
Simplest $v_2^q = v_2^s$ case: $v_2^{\text{hadron}}(np_\perp)/n$ universal fn. ($n = 2, 3$ for meson, baryon)

Experimental test of v_2 scaling

Simplest $v_2^q = v_2^s$ case: $v_2^{hadron}(np_{\perp})/n$ universal fn. ($n = 2, 3$ for meson, baryon)

Castillo [STAR], HIC03: K_S^0, Λ, Ξ flow

PHENIX, PRL91 ('03): π, K, p



- **coalescence predictions confirmed** for $\pi, K, K_0, p, \Lambda, \Xi$ - also for Ω, ϕ (poor statistics)
- pions are little off the curve, likely due to resonance decays [Greco & Ko]
- **surprisingly, RHIC data indicate $v_2^q \approx v_2^s$** \rightarrow next open question: **charm v_2 ?**

Wave function

Hadron rest frame: - consider $|\vec{p}_1, \vec{p}_2\rangle = |\vec{q}/2, -\vec{q}/2\rangle$ (meson)
- with $\langle |\vec{q}| \rangle \sim \Lambda_{QCD}$

Fast, weakly bound hadron w/ momentum $\vec{n}p$: - a kinematic estimate

$$p'_i \equiv \vec{p}'_i \cdot \vec{n} = \frac{E'_i}{m_H} p + \vec{q}_i \cdot \vec{n} \frac{\sqrt{p^2 + m_H^2}}{m_H} \approx \frac{m_i}{m_H} p + \frac{\vec{q}_i \cdot \vec{n} p}{m_H}$$

→ with momentum fractions $z_i \equiv p'_i/p$

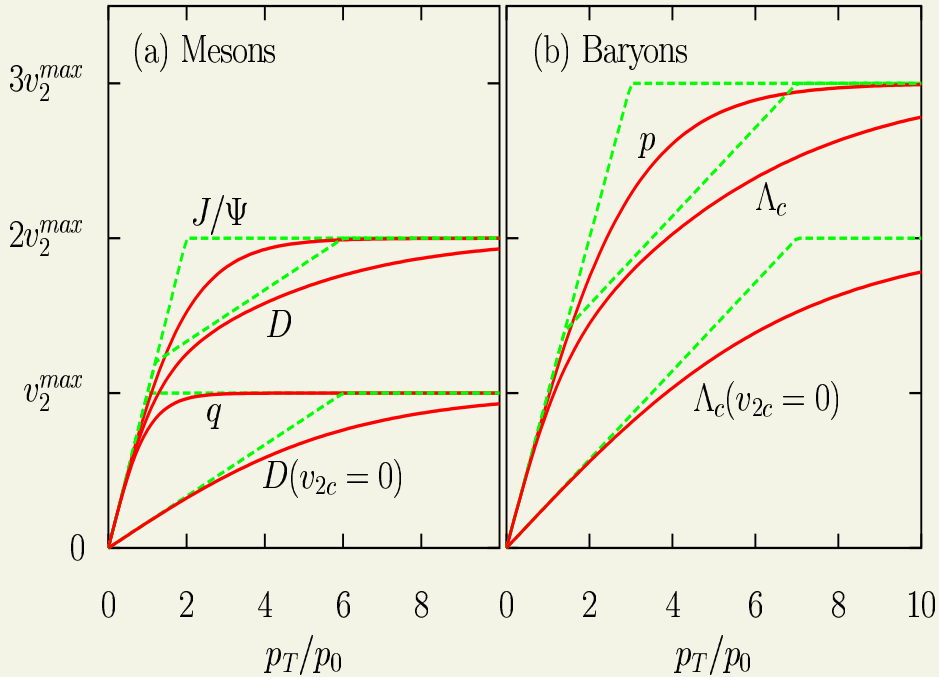
$$z_i = \bar{z}_i + \delta z_i \approx \underbrace{\frac{m_i}{m_H}}_{\text{average}} + \underbrace{\frac{\vec{q}_i \cdot \vec{n}}{m_H}}_{\text{spread}}$$

- Two effects:**
- **lighter quark carries smaller fraction of momentum**
- because similar velocities (not momenta) in moving frame
 - **wave function is narrower in δz for heavy quarks**

Elliptic flow for charm

generic expectations:

Lin & D.M. ('03)

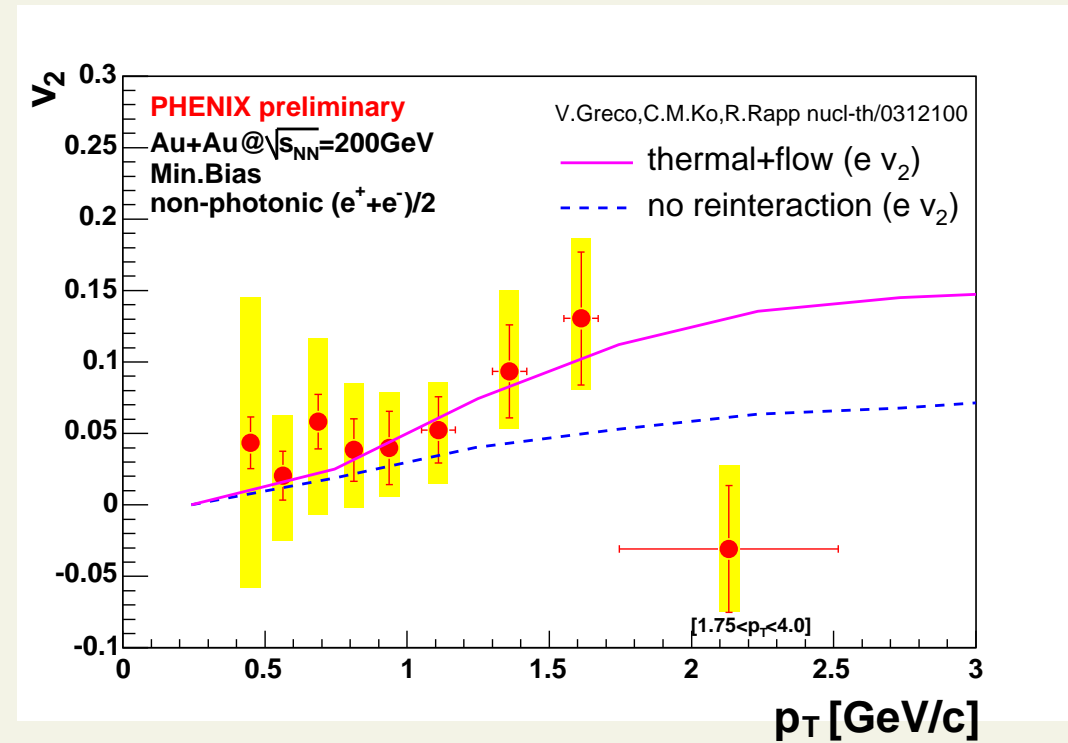


simplicistic (linear & flat); more realistic (based on parton cascade MPC)

- $v_2(p_\perp)$ rises slower, saturates later for asymmetric systems (D, D_s, Λ_c)
 - heavy quark carries most of the hadron momentum (momentum proportional to constituent mass)
- nonzero v_2 for D, D_s, Λ_c even for zero charm v_2 (no thermalization)

decay electron v_2 : $D \rightarrow K \nu e, D \rightarrow K^* \nu e$

Kaneta, nucl-ex/0404014; Greco & Ko, nucl-th/0312100:



e's from hadron decays and γ -conversion subtracted
 \equiv "non-photonic"

- need better data to discriminate between scenarios with or without charm flow

Coal. formalism for diffuse freezeout

Gyulassy, Frankel & Remler (GFR): [NPA 402, 596 ('83)]

- for each constituent pair/triplet, propagate particles to the **latest of freezeout times** and evaluate weight $W(\Delta x, \Delta p)$ there
- reason: **any** interaction would break up a weak bound state
- note, **relative distance changes(!)**, e.g., **if** $t_2 > t_1$:

$$weight = W_M(\vec{x}_1(t_1) + (t_2 - t_1)\vec{v}_1 - \vec{x}_2(t_2), \vec{p}_1 - \vec{p}_2)$$

- Naturally incorporates:**
- “diffuse” 4D freezeout
 - space-time and space-momentum correlations

Model ingredients

Processes: ideally: $f - 2 \rightarrow 2$ parton scatterings, showers ($1 \rightarrow 2, 1 \rightarrow 3$), parton fusion ($2 \rightarrow 1, 3 \rightarrow 1$), inelastic $n \rightarrow m$, parton recombination to hadrons, hadron breakup, ... etc.

here:

- only $2 \rightarrow 2$ (with $g, u, d, s, c, \bar{u}, \bar{d}, \bar{s}, \bar{c}$)
- no parton showers until freezeout
- coalescence rate computed over freezeout 4D volume via GFR
- partons with no coalescence partner fragment as in vacuum

Coalescence part:

- assume easy color neutralization - no color penalty factors
- but consider spin & flavor
- channels: $\pi, K, \eta, \eta'; \rho, K^*, \omega, \Phi, D, D^*, J/\psi, \eta_c$
 $p, n, \Sigma, \Lambda, \Xi; \Delta, \Omega$
- “spherical box” Wigner functions:
$$W_M = \Theta(p_M - |\Delta p|)\Theta(x_M - |\Delta x|)$$
$$W_B = \prod_{k \neq i, j} \Theta(p_B - |\Delta p_{ij}|)\Theta(x_B - |\Delta x_{ij}|)$$
$$x_M = x_B = 1 \text{ fm}$$
- convert g to a random q (extreme case of $q - \bar{q}$ splitting)
- when several coalescence final states, unbiased random choice of one

Codes:

- MPC 1.8.0 for parton transport
- JETSET 7.4.10 for fragmentation & decays (“out of box”)

Various issues regarding coalescence

- **coalescence dynamics**
- **extend to low- p_T regime:** **simple formula not applicable** (unitarity)
 - does coal w/ unitarity still reduce to **statistical models/hydrodynamics?**
- **hadron correlations:** **coal dilutes jetlike correlations??**
 - **not necessarily:** constituent correlations possible \Rightarrow map uniquely to hadron correlations
 - **a correlation source:** **parton showers** [Hwa & Yang, nucl-th/0401001]
- **binding energy:** only relevant for **Goldstones π & K**
 - likely a **small correction** because resonance channels $K^* \rightarrow K; \rho, \omega, \Delta, \dots \rightarrow \pi$ dominant
- **constituents vs partons:** so far - two, mostly independent components
 - **Hwa & Yang:** hadrons made of constituents, which in turn consist of partons
limited to momentum space \rightarrow need **extension to 6D phasespace**
- **entropy:** **$2 \rightarrow 1, 3 \rightarrow 1$??**
 - **globally, problem is cured** by resonance decays and larger s/n for massive particles